The Frenet formulas for a Weyl space

By

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Blaschke recently gave (¹) the Frenet formulas for a curve that is traced in a space (R_n) with a Riemannian metric in which one defines parallel displacement as Levi-Cività did (²). Upon employing the same calculation procedures that Blaschke did, we have obtained the *Frenet* formulas for a curve that is traced in a space (W_n) with a Weyl metric.

An *n*-dimensional Weyl (³) space (W_n) is an *n*-dimensional multiplicity in which the metric is defined by two forms (one quadratic and one linear):

$$ds^2 = \sum_{i,k=1}^n g_{ik} dx_i dx_k , \qquad \qquad d\varphi = \sum_{i=1}^n \varphi_i dx_i .$$

 $d\varphi$ is an invariant for any continuous transformation (T) of the form:

$$x_i = \psi_i (y_1, ..., y_n)$$
 $(i = 1, 2, ..., n)$

Moreover, if one makes a change of calibration – i.e., if one takes a unit of length that is $\sqrt{\lambda}$ times smaller (λ = continuous function of $x_1, ..., x_n$) – then the two forms will become:

$$ds'^2 = \sum_{i,k=1}^n \lambda g_{ik} dx_i dx_k$$
, $d\varphi' = d\varphi - \frac{d\lambda}{\lambda}$

The laws of geometry must be satisfied under the following two conditions:

1. They are expressed by formulas that are invariant under any transformation (T).

^{(&}lt;sup>1</sup>) Math. Zeit. **6** (1919).

⁽²⁾ Rendiconti del Circolo mat. di Palermo **42** (1917).

^{(&}lt;sup>3</sup>) See WEYL, Raum, Zeit, Materie, 4th edition, § 16.

2. Those formulas remain invariant if one changes g_{ik} into λg_{ik} and φ_i into $\varphi_i - \frac{1}{\lambda} \frac{\partial \lambda}{\partial x}$.

Weyl define parallel displacement in relation to this new concept. Let a vector with components $(\xi^1, ..., \xi^n)$ be attached to the point $P(x_1, ..., x_n)$. We say that its *measure* is:

$$m = \sum_{i,k=1}^{n} g_{ik} \, \xi^i \xi^k$$

Upon displacing *P* to $P'(x_i + dx_i)$ by congruence, its components will become $\xi^i + d\xi^i$, with:

$$d\xi^{i} = -\frac{1}{2} \sum_{k,r,t=1}^{n} g^{ik} \left[\frac{\partial g_{ik}}{\partial x_{r}} + \frac{\partial g_{kr}}{\partial x_{t}} - \frac{\partial g_{rt}}{\partial x_{k}} + g_{rk}\varphi_{t} + g_{tk}\varphi_{r} - g_{rt}\varphi_{k} \right] \xi^{r} dx_{t}$$

Let C be a curve whose parametric equations are $x_i = x_i$ (s). Imagine that we have fixed a vector Ξ at each point *P* (*s*) whose components (ξ^i) are continuous functions of *s* according to a continuous law. Let P(s) and P'(s + ds) be two neighboring points, so they will then correspond to the two vectors Ξ and Ξ' . Displace Ξ from P to P'by congruence, so one will get a vector Ξ^* at P that is generally different from Ξ' . The difference $\Xi' - \Xi^*$ is an infinitely small vector that is attached to the arc PP'. Form:

$$\theta\left(\Xi\right)=\frac{\Xi'-\Xi^*}{ds}.$$

We will then get a new vector $\theta(\Xi)$ that is attached to the point P(s) of the curve C and depends upon the field Ξ in an invariant manner (¹). Then let $\Xi = \Xi_1$ with the components $\xi_{(1)}^i = dx_i / ds$. Then set:

$\theta(\Xi_1) = \Xi_2$	with components	$\xi^{i}_{(2)},$
$\theta(\Xi_2) = \Xi_3$	" "	$\xi^{i}_{(3)},$
$\theta(\Xi_{n-1}) = \Xi_n$	with components	$\xi_{(n)}^i$.

The *n*-hedron $\Xi_1, \Xi_2, ..., \Xi_n$ is not orthogonal, in general. Orthogonalize it using Schmidt's method $\binom{2}{2}$ by defining an *n*-hedron that is composed of the *n* vectors:

^{(&}lt;sup>1</sup>) WEYL, *loc. cit.*, pp. 103.
(²) "Integralgleichungen, etc....," Math. Ann. 63.

$$H_{p} = \begin{vmatrix} (1,1) & (1,2) & \cdots & (1,p-1) & \Xi_{1} \\ (2,1) & \cdots & \cdots & \vdots & \vdots \\ \vdots & \vdots & \vdots & \vdots & \vdots & \vdots \\ (p,1) & \cdots & \cdots & (p,p-1) & \Xi_{p} \end{vmatrix} \qquad (p = 1, 2, ..., n),$$

in which one sets:

$$(\alpha, \beta) = \sum_{i,k=1}^{n} g_{ik} \xi_{(\alpha)}^{i} \xi_{(\beta)}^{k}, \qquad D_{p} = \begin{vmatrix} (1,1) & \cdots & (1,p) \\ \vdots & \ddots & \vdots \\ (p,1) & \cdots & (p,p) \end{vmatrix}, \qquad D_{0} = 1.$$

The *n*-hedron (*N*), $H_1, ..., H_n$ is orthogonal and normalized; i.e.:

$$\sum_{i,k=1}^{n} g_{ik} \eta^{i}_{(\alpha)} \eta^{k}_{(\beta)} = \delta_{pq} = \begin{cases} 1 & \text{is} \quad p = q \\ 0 & \text{is} \quad p \neq q. \end{cases}$$

The Frenet formulas for the curve C are the formulas that give the values of s :

$$\theta\left(H_p\right) = \frac{W_p' - W_p}{ds} \,.$$

One finds, by some simple calculations, that:

$$(F) \begin{cases} \theta \eta_{(1)}^{i} = \frac{1}{2} \frac{d\varphi}{ds} \eta_{(1)}^{i} + \frac{1}{\rho_{1}} \eta_{(2)}^{i}, \\ \theta \eta_{(2)}^{i} = -\frac{1}{\rho_{1}} \eta_{(1)}^{i} + \frac{1}{2} \frac{d\varphi}{ds} \eta_{(2)}^{i} + \frac{1}{\rho_{2}} \eta_{(3)}^{i}, \\ \dots \\ \theta \eta_{(p)}^{i} = -\frac{1}{\rho_{p-1}} \eta_{(p-1)}^{i} + \frac{1}{2} \frac{d\varphi}{ds} \eta_{(p)}^{i} + \frac{1}{\rho_{p}} \eta_{(p+1)}^{i}, \\ \theta \eta_{(p)}^{i} = -\frac{1}{\rho_{n-1}} \eta_{(n-1)}^{i} + \frac{1}{2} \frac{d\varphi}{ds} \eta_{(n)}^{i}, \\ \theta \eta_{(n)}^{i} = -\frac{1}{\rho_{n-1}} \eta_{(n-1)}^{i} + \frac{1}{2} \frac{d\varphi}{ds} \eta_{(n)}^{i}, \end{cases}$$
in which:

$$\rho_k = \frac{D_k}{\sqrt{D_{k-1}D_{k+1}}}.$$

The $\theta \eta_{(k)}^i$ are then homogeneous linear functions of the $\eta_{(q)}^i$. The determinant of those functions is *skew-symmetric*; the ρ_i are the (n-1) radii of curvature of the curve. That determinant

possesses a principal diagonal whose terms are all equal to $\frac{1}{2} \frac{d\varphi}{ds}$. For a space (R_n), the formulas (F) will be the same as the ones that we just found, except that all of the terms in the principal diagonal will be equal to zero. (The ρ_i will not have the same value, since they depend upon the φ_i .) If one regards the trihedron (N) as moving along the curve C then one can say that one passes from one of its position to the neighboring position by displacing by congruence and then subjecting it to a rotation that is defined by the curvatures $1 / \rho_i$ of C, and finally deforming it with a homothety of ratio $1 + d\varphi/2$.